Physically Unclonable Cryptographic Primitives by Chemical Vapor Deposition of Layered MoS$_2$

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Supporting Information

Abstract: Physically unclonable cryptographic primitives are promising for securing the rapidly growing number of electronic devices. Here, we introduce physically unclonable primitives from layered molybdenum disulfide (MoS$_2$) by leveraging the natural randomness of their island growth during chemical vapor deposition (CVD). We synthesize a MoS$_2$ monolayer film covered with speckles of multilayer islands, where the growth process is engineered for an optimal speckle density. Using the Clark–Evans test, we confirm that the distribution of islands on the film exhibits complete spatial randomness, hence indicating the growth of multilayer speckles is a spatial Poisson process. Such a property is highly desirable for constructing unpredictable cryptographic primitives. The security primitive is an array of 2048 pixels fabricated from this film. The complex structure of the pixels makes the physical duplication of the array impossible (i.e., physically unclonable). A unique optical response is generated by applying an optical stimulus to the structure. The basis for this unique response is the dependence of the photoemission on the number of MoS$_2$ layers, which by design is random throughout the film. Using a threshold value for the photoemission, we convert the optical response into binary cryptographic keys. We show that the proper selection of this threshold is crucial for maximizing combination randomness and that the optimal value of the threshold is linked directly to the growth process. This study reveals an opportunity for generating robust and versatile security primitives from layered transition metal dichalcogenides.

Keywords: security, cryptographic primitives, physically unclonable, MoS$_2$, CVD growth

Modern society demands information security.\(^1\) Globalization of supply chains has undermined trust in electronic devices, which were once manufactured entirely by a single trusted factory. Further, the ubiquity of today’s advanced manufacturing poses additional challenges, because such resources are now more accessible to adversaries for developing sophisticated security attacks. A vast number of such attacks are physical\(^2,3\) and range from counterfeiting to unauthorized access. As a result, authentication of electronic devices and information has become increasingly important.

Physically unclonable functions (PUFs) are among promising security primitives for entity identification or cryptographic key generation.\(^4,5\) A variety of PUFs have used manufacturing variability\(^6,7\) or materials disorders\(^8–13\) for generating a security key on demand. Specifically, applying a challenge (such as an electrical or an optical stimulus) to a PUF produces a unique response (a security key). Hence, this concept generates security keys that are unique for each electronic device. A PUF construct must be easy to produce at low cost and yet have a physical structure that is impossible to replicate, i.e., physically unclonable. Notably, the uniqueness of the response and the physical unclonability are the core defining properties of a PUF.\(^5\)

Silicon-based PUFs have garnered significant attention as a potential low-cost solution for securing digital systems, due to their compatibility with the complementary metal-oxide-semiconductor (CMOS) technology.\(^7\) These PUFs use the intrinsic variability of the CMOS process, which follows a Gaussian distribution with a mean value close to zero.\(^14\) Because of this attribute of standard CMOS processes, the measured output of all silicon PUFs requires extensive data processing, which aims to enhance the overall robustness and uniqueness of the PUF response.\(^5,6,15,16\) Examples of postprocessing procedures include error correction, masking of unreliable bits, addition of helper data, and fuzzy extraction. Not only does the extensive postprocessing add to the complexity of silicon PUFs, but it might also lead to vulnerability to physical attacks.\(^17\) Hence, there is an
opportunity for realizing nonsilicon cryptographic primitives that can generate inherently more unique and robust security keys. A large body of work has reported the natural randomness of nanomaterials, which can potentially be leveraged for constructing nonsilicon security primitives. Optimizing the security and reliability of these primitives requires tuning the conditions of the physical process involving nanomaterials (e.g., synthesis or assembly). Further, the scalability of the physical process is important for producing primitives on a large scale and thus realizing a viable nanomaterials-based security technology. A recent work of Hu et al. provides a meaningful experimental demonstration of this concept, implementing carbon nanotube (CNT) random bits by leveraging the randomness of a CNT assembly process. Beyond CNTs, the prospects of two-dimensional (2D) transition metal dichalcogenides (TMDs) for security applications are still unexplored.

Several key advances in the field of 2D materials make this exploration worthwhile. First, many studies have reported the large-scale growth of various TMDs on amorphous substrates with apparent random nucleation. A true island growth is known to exhibit complete spatial randomness (CSR), which indicates that the occurrences of nucleations are independent and equally probable everywhere on the substrate. On the basis of these properties, the natural randomness of a true island growth offers a close approximation of true randomness. Hence, a TMD film with CSR nucleations is appealing for constructing highly unpredictable security primitives. However, until now, no study has statistically tested whether TMD island growth is CSR. The second advance relates to discovering the strong thickness dependence of photoemission from most semiconducting TMDs. The illustration in Figure 1a shows this concept. Combining this unique property with the randomness of the TMD growth can result in a PUF that provides a random response to an optical stimulus. Finally, the substrate-agnostic property of these 2D materials makes the resulting security primitives versatile by allowing their integration with different platforms from conventional digital systems to emerging flexible electronics.

Here, we introduce a physically unclonable security primitive constructed from layered molybdenum disulfide (MoS2). We used MoS2 as the model system since it has been studied heavily in the family of 2D layered TMDs. Figure 1 illustrates the proposed concept. A large-area MoS2 film is produced using a chemical vapor deposition (CVD) process in a layer-plus-island growth mode (Figure 1b). The details of all CVD growth experiments are given in Supporting Information. Figure 1c shows the photo of a CVD MoS2 sample indicating two distinct regions of growth. Scale bar is 5 mm. An array of 32 x 64 pixels was then formed in region II. (e) Stimulating the physical MoS2 primitive with a laser light produces a unique optical response with randomly distributed ON and OFF pixels.

Figure 1. MoS2-based physically unclonable security primitives: (a) Schematic illustration of the energy band structure of monolayer and multilayer MoS2, indicating strong dependence of excitonic emission on the number of MoS2 layers. (b) CVD using solid-phase precursors was used for synthesis of large-area MoS2 films. (c) Photo of a CVD MoS2 sample indicating two distinct regions of growth. Scale bar is 5 mm. (d) An array of 32 x 64 pixels was then formed in region II. (e) Stimulating the physical MoS2 primitive with a laser light produces a unique optical response with randomly distributed ON and OFF pixels.
tests, we confirm the randomness and stability of the security keys generated from the proposed security primitive.

RESULTS AND DISCUSSION

Layer-Plus-Island Growth of CVD MoS2. Among the different methods used for growing TMDs, CVD techniques offer better control over thickness on a large scale.20,32−34 We grew large-area MoS2 films onto 285 nm SiO2 on p+ silicon substrates by CVD from sulfur and MoO3 precursors.35 Figure 1b schematically illustrates the CVD reactor based on the solid-phase precursors. Two distinct growth regions are typically evident along the substrate (see Figure 1c), indicating that in this CVD process the growth depends on the distance of the substrate from the MoO3 powder. Region I is the farthest from the MoO3 powder in the reactor, where the optimal growth conditions yield a continuous monolayer. Region II, located closer to the MoO3 powder, is covered by a continuous monolayer film with randomly distributed multilayer islands. According to the surface science of thin film growth, the growth mode strongly depends on the deposition rate of the growth species and the substrate temperature.24 It is known that the growth mode will deviate from the layer-by-layer mode to the layer-plus-island mode once the deposition rate exceeds a critical value. This explains the presence of these two prominent growth modes along the substrate.20,24 Indeed, the layer-by-layer growth occurs in region I with low Mo vapor pressure and the growth follows a site-saturated growth kinetics (see Supporting Information). Despite the spatial randomness of the nucleation sites, region I of our samples is suboptimal for constructing a dense array of random binary code because of the relatively sparse spatial distribution of the multilayer films grown mostly at the grain boundaries (spacing from 20 to 80 μm). In contrast, region II (closer to the MoO3 powder) is exposed to a higher concentration of Mo vapor, resulting in the layer-plus-island growth mode and thus the random nucleation of multilayer islands on the monolayer film, as shown in Figure 2a. To produce the physically unclonable security primitive, we engineer the growth process in this region to achieve an optimal surface coverage of the multilayer islands.

Engineering and Testing the Growth Randomness. Complete spatial randomness is central to constructing strong cryptographic keys in our proposed concept. To test for CSR of the island growth in region II, we apply the statistical test by Clark and Evans31 on images from this region, taken at an early stage of the multilayer nucleation on the continuous monolayer film, e.g., Figure 2a at t = 60 s. If the island growth is CSR, then the distribution of nearest neighbor distances (i.e., the distances between the islands and their nearest neighbor) has a mean $r_{CSR} = 1/(2\sqrt{\rho})$ and a variance $\sigma^2_{CSR} = (4 - \pi)/(4\pi\rho)$, where $\rho$ is the particle density per unit area. Therefore, we can test for CSR by testing the null hypothesis that the mean of nearest neighbor distances is equal to $r_{CSR}$. Using the two-tailed test for the population mean, we compute the standard Z-score given by

$$Z = \frac{\langle r \rangle - \langle r_{CSR} \rangle}{\sqrt{\sigma^2_{CSR}/N}}$$  \hspace{1cm} (1)$$

where $r_i$ is the sample mean of the nearest neighbor distances computed with $N$ particles. That is,

$$r_i = \sum_{j=1}^{N} r_{ij} / N$$  \hspace{1cm} (2)$$

where $r_{ij}$ is the nearest neighbor distance of the $i$th island. At a 0.05 significance level, the null hypothesis is to be rejected if $Z$...
The hypothesis that the mean nearest neighbor distance is at the 0.05-level of significance, we cannot reject the null hypothesis that the mean nearest neighbor distance is $r_{CSR}$. This suggests that the island growth in region II exhibits CSR, hence all nucleations are independent and the probability of nucleation is the same everywhere on the surface.

Considering CSR island growth in region II, the Avrami equation can be used to draw insight into the time evolution of the island growth. For a growth time $t$, the fractional surface coverage $f$ of the multilayer islands ($\geq$2L) is approximated by

$$f_{\geq 2L}(t) = 1 - e^{-kt^n}$$  \hspace{1cm} (3)

where the Avrami exponent $n$ gives information about the kinetics of the island growth. To analyze the growth kinetics, we prepared several samples strictly by varying the growth conditions, including the quantity of the precursors, the sample size, and the sample position relative to MoO$_3$. We then imaged the samples to compute the fractional areal coverage of monolayer and multilayer films in region II, as shown in Figure 2a. Assuming time-invariant growth kinetics, this experiment provides a good approximation of the time evolution of the surface coverage.

From the optical images, we made two key observations. First, the nucleation is continuous from the concurrent presence of thin (mostly bilayer) and thick islands in all different stages of the growth. Second, the growth is mostly 2D, i.e., the lateral dimensions of islands grow faster than the thickness. Figure 2b summarizes the time evolution of the fractional surface coverage for the monolayer film and the multilayer islands. In this plot, $t = 0$ represents the time at which a monolayer film fully covers the surface in Region II.

From our observations, this time corresponds to the beginning of the growth cycle at 850 °C. From the data, the multilayer islands cover 50% of the monolayer film surface at $t \approx 240$ s. In Figure 2c, we plotted $\ln(-\ln(1 - f_{2L}))$ as a function of $\ln(t)$, where the slope of the fitted line gives the estimate for the Avrami exponent $n$. We found $n \approx 2$ for our growth experiments, suggesting a 2D disk-shaped growth governed by the surface diffusion. Eq 3 provides a reasonable fit to the experimental data in Figure 2b, further confirming CSR of the multilayer island growth in region II. In addition, the inflection point of the fitted curve at $t \approx 180$ s corresponds to the crossover from the isolated island growth to the island overlap growth.

After analyzing the growth kinetics, we adjusted the growth time to obtain MoS$_2$ films with equal surface areas of exposed monolayer (1L) film and of the multilayer islands, i.e., $f_{1L} = f_{\geq 2L} = 0.5$. This was done to achieve the maximum combination randomness in the security key responses.

**Cryptographic Key Generation.** To implement the MoS$_2$-based security primitives, we fabricated dense arrays consisting of 32 × 64 pixels from the film in region II (see Methods). These arrays have a pixel size of 2 μm × 2 μm and an equal pixel spacing of 2 μm. This pixel size was chosen because it is comparable with the dimensions of the state-of-the-art CMOS image sensors. Figure 3a shows the optical image of a 2D MoS$_2$ array with 2048 pixels, fabricated on a SiO$_2$/Si substrate. Due to the randomness of the multilayer island growth on a continuous monolayer film, the content of each pixel is random. A pixel might consist of a monolayer, a multilayer, or a mixture of the two. Figure 3b illustrates the zoomed-in view of three neighboring pixels. These pixels...
the thickness dependence of the excitonic emissions in MoS2.


Figure 4. Optical response and security metrics of the MoS2 primitive: (a) Stimulating the 2D MoS2 array with a laser light produces an optical response that is unique to this primitive. (b) The photoemission spatial map was converted to a 2D binary array by comparing each pixel with the ON/OFF threshold. (c) Standard security tests confirm uniqueness and repeatability of the security keys. (d) By studying the aging properties of the photoemission for multiple MoS2 pixels, we confirm that the MoS2 primitives are highly stable. The red dashed line is guide to the eye and has a slope of 1.

After fabrication, the physical security primitive was stimulated using a laser light to generate an optical response. We expect the response to be unique to the security primitive given the random thickness distribution of the CVD MoS2 and the thickness dependence of the excitonic emissions in MoS2. Figure 3c confirms the material type (which is MoS2 here) and the corresponding thickness, determined from the distance between the peak position of the in-plane ($E_{2g}$) and the out-of-plane ($A_{1g}$) phonon modes.

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We have previously shown that the growth process was tuned to obtain equal surface coverage of 50% (see Figures 3e) represents the most ambiguous case for classifying a pixel as ON or OFF. Therefore, we use the photoemission of such a mixed pixel as the photoemission threshold $\theta$.

The data in Figure 3e shows clearly that the ratio of the PL intensity of the mixed pixel to that of a full monolayer pixel (about 0.12) is noticeably smaller than the monolayer areal coverage of the mixed pixel (0.5). We noticed that this trend is consistent across the different arrays on the sample. To better understand this observation, we closely examined the PL spectra of multiple pixels with varying monolayer fractional surface coverage, from 0 (full bilayer) to 1 (full monolayer). We selected these pixels at random from different locations within the same array. To accurately determine the surface coverage of the monolayer film within each pixel, we performed atomic force microscopy (AFM). The corresponding AFM images and the PL spectra of these pixels are shown in Supporting Information. We calculated the total area under the PL emission curve for each pixel in the wavelength range of 580–770 nm (the integrated photoemission). Then, the integrated PL of the pixels were normalized with respect to that of a full monolayer pixel. To account for the difference in the monolayer content of these pixels, we divided the normalized integrated photoemission of each pixel by its monolayer fractional coverage. These final values provide a close approximation of the quantum yield (QY).

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choice on the behavior of the security primitive. Considering CSR of the island growth and the equal surface coverage by the monolayer and multilayer MoS$_2$, it is expected that the distribution of the ON and OFF pixels shows no or weak dependency on the pixel size and the pixel spacing in the 2D MoS$_2$ array. We confirmed this by fabricating multiple arrays with different pixel sizes and spacings, where the arrays demonstrate equal distribution of random ON and OFF pixels (Supporting Information). Hence, the strength of the security primitive is robust to the pixel choice and spacing choice.

Analyzing Security and Stability Metrics. We next analyzed the security metrics of the 2D binary array. Three important metrics are typically used to evaluate the strength of a security primitive: uniqueness, repeatability, and uniformity.

Uniqueness is the ability of a key to be distinguished from other keys. We use the average Hamming interdistance to quantify uniqueness. The Hamming interdistance between two keys is the minimum number of bit substitutions required to transform one key to another. The 32 rows of the 2D binary array are 64-bit security keys to be tested. We compute the Hamming interdistance of all 496 possible pairs of keys (see Supporting Information). Figure 4c shows the Hamming interdistance distribution. A binomial distribution with parameters $p = 0.495$ and $N = 64$ provides a good fit based on the Kolmogorov–Smirnov test. The inverse of the binomial distribution at cumulative probability 0.05 is 25. This means that for two randomly generated 64-bit keys, there is a 95% probability that the keys differ in at least 25 bits. Hence, there is a 95% chance that it will require at least 64 choose 25 ($\binom{64}{25}$) worst-case number of attempts to guess an unknown key from another known key.

A random key must also produce a consistent response to a given input challenge. The difference in response of a given binary key to the same challenge is quantified by the Hamming intradistance, which represents the repeatability of the random binary code. Therefore, the ideal intradistance is zero. Figure 4c shows the results of the Hamming intradistance, indicating high repeatability of the MoS$_2$ security keys. The observed bit error rates are measurement artifacts and originate from the limited spatial accuracy of the automated sample stage of the PL measurement system. To investigate the long-term stability of our security primitives, we measured a random sample of 200 pixels after 6 months storage in ambient air, as shown in Figure 4d. In this experiment, the candidate pixels had either high fractional areal coverage by a monolayer or a multilayer MoS$_2$ film. The unchanged emission properties of these pixels indicate that our MoS$_2$ cryptographic primitives are highly stable.

To maximize the combination randomness of a binary array, each pixel should have an equal probability (i.e., 0.5) of being ON or OFF. This is defined as the uniformity property and is quantified by the Hamming weight of the key. Specifically, the Hamming weight indicates the number of bit substitutions to convert the key to an array of all zeros and has an ideal value of 0.5. We calculate the normalized Hamming weight on all 32 64-bit rows of the 2D binary array, and found the average to be 0.48 (see Supporting Information). As described earlier, the measured Hamming weight is directly linked and controlled by selection of the photoemission threshold. This discussion explains a key design parameter, which must be determined properly to achieve maximum combination randomness. Through our observations of the CVD growth process, we can tune the photoemission threshold value, thus taking full advantage of the inherent randomness of the cryptographic primitive.

Lastly, we comment on the physical unclonability of our security primitive. To classify a random physical construct as a security primitive, it must be capable of preventing an unauthorized duplication of its physical structure, even if the adversary has full control of the PUF manufacturing. Figure 5 illustrates the topographic images of several MoS$_2$ pixels from the 2D array. These images highlight the extreme intricacy of the pixel structures. The heterogeneity of the thickness and the randomness of the spatial distribution of MoS$_2$ films within the pixels—originating from CSR of the multilayer island growth—make the 2D MoS$_2$ array physically unclonable. In contrast, we have shown that a physical array of pixels constructed from randomly distributed monolayer islands in Region I (before they coalesce) can be replicated using standard nanofabrication techniques (see Supporting Information). This experiment underscores that the mere randomness of a physical construct, although necessary, is insufficient for realizing a security primitive.

CONCLUSIONS

We introduce a physically unclonable security primitive based on layered MoS$_2$. Two fundamental properties underlie this security technology: (i) CSR of multilayer island growth during CVD of MoS$_2$, confirmed by Clark–Evans test, and (ii) strong thickness dependence of photoemission in MoS$_2$. Using the Avrami equation, we show that the nucleation is continuous and that the island growth is two-dimensional. We define a photoemission threshold to convert the optical response of the physical 2D array into a binary cryptographic key. Our experiments establish that this threshold is directly linked to the engineered growth process and can be tuned to maximize the combination randomness of the security keys. These security primitives are easy to produce on a large scale using CVD and yet impossible to duplicate because of the intricate physical structures of the pixels. The findings of this study can be readily extended for the development of physically unclonable primitives based on other semiconducting transition metal dichalcogenides.
METHODS
We performed CVD growth using MoO₃ and sulfur solid precursors without requiring a growth promoter. The growth was performed using a custom-made setup at 850 °C with a nitrogen flow of 10 sccm. The optimal quantities of MoO₃ and sulfur precursors are about 6 mg and 100 mg. In all the experiments, the films were grown in the presence of excess sulfur. The details of the growth experiments are given in the Supporting Information. The 2D MoS₂ array was fabricated using an e-beam lithography step followed by patterning in an CF₆/O₂ plasma. The 2D arrays were stimulated using a green laser for producing an optical response.

ASSOCIATED CONTENT
Supporting Information
The Supporting Information is available free of charge on the ACS Publications website at DOI: 10.1021/acsnano.7b07568.
Details of chemical vapor deposition, analysis of growth in Regions I and II, interpretation of Avrami exponent, normalized quantum yield analysis, physical unclonability experiments, details of PL measurements, studies of pixel size and spacing choices, and analysis of security metrics (PDF)

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The authors declare no competing financial interest.

ACKNOWLEDGMENTS
This work was supported in part by NSF Award #1638598 and by U.S. Army Research Office Award #W911NF-16-1-0301. This research used resources of the Center for Functional Nanomaterials, which is a U.S. DOE Office of Science Facility, at Brookhaven National Laboratory under Contract No. DE-SC0012704. The authors acknowledge C. Black and J. Uichanco for helpful discussions.

REFERENCES


